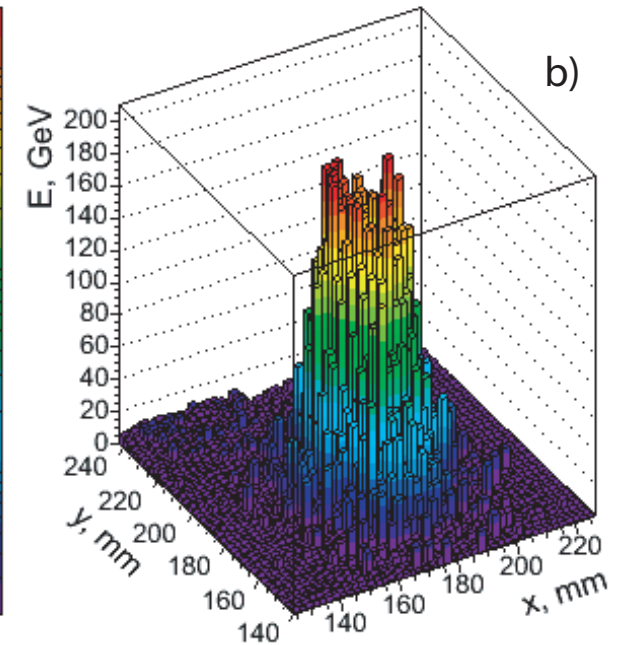
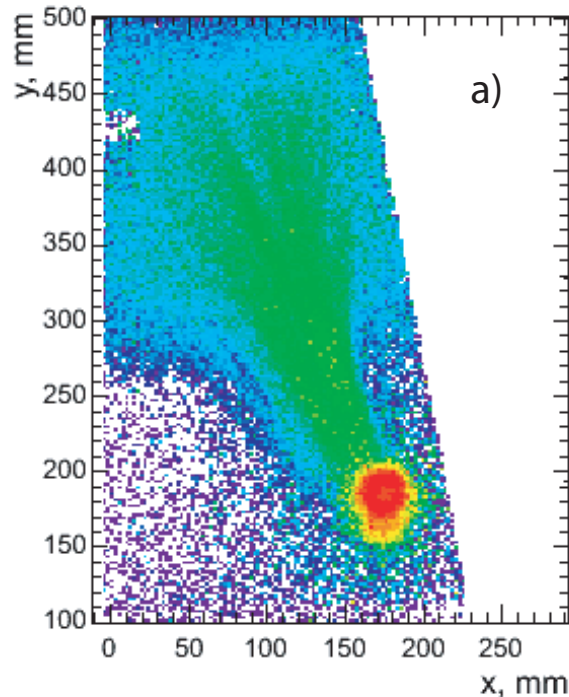


An AntiCherenkov Photomultiplier

Nural Akchurin, Jordan Damgov, Sung-Won Lee and Efe Yazgan
Texas Tech University

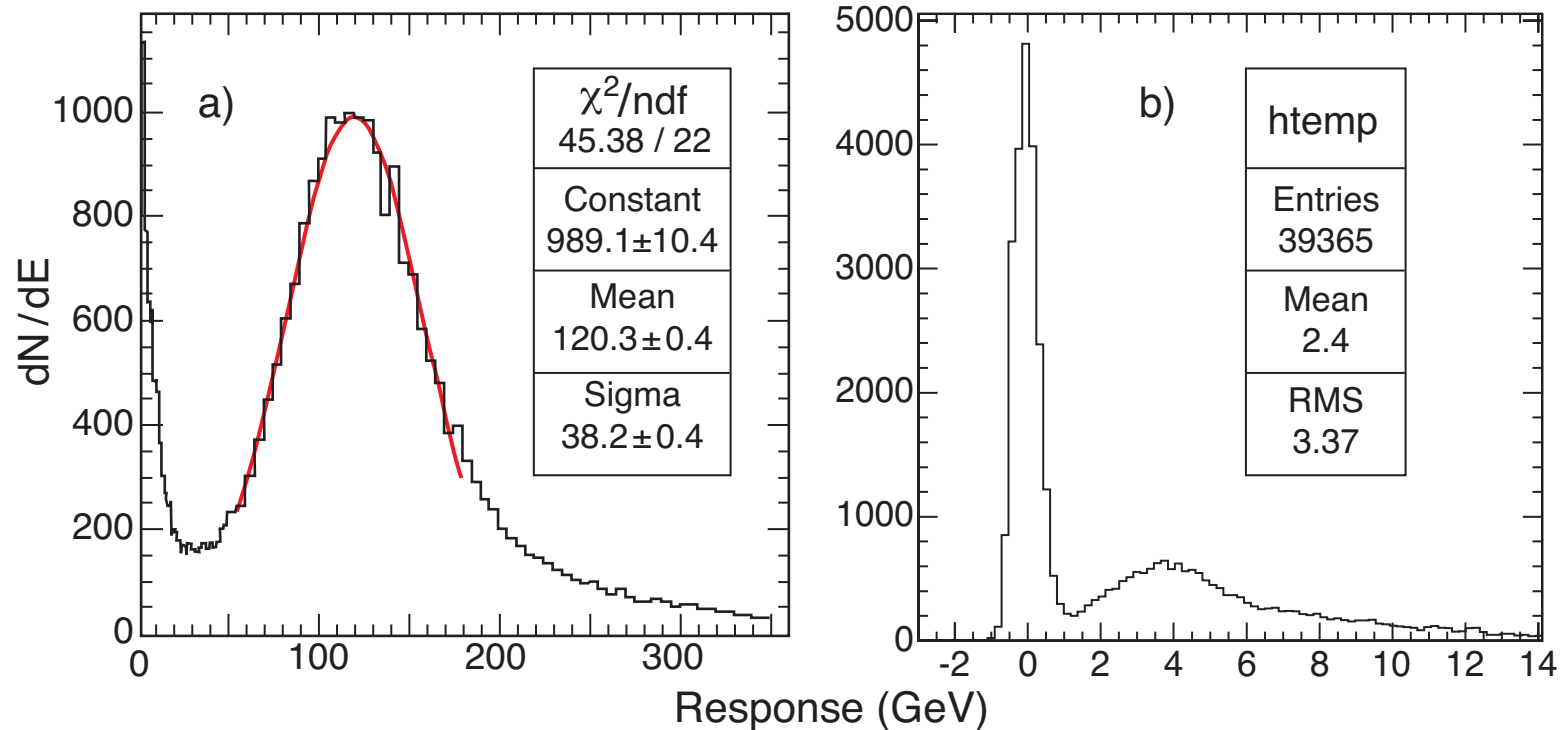
Problem?



“Hot spots” have been observed in many detectors, especially in calorimeters, for when the signal strength due to instrumentation exceeds that of the desired signal. This becomes a concern when the light yield of a calorimeter is low (a few pe per GeV) and there is no redundancy in the system to veto these events

Problem ?

S. Abdullin *et al*, EPJC, **53**, 139-166 (2008)

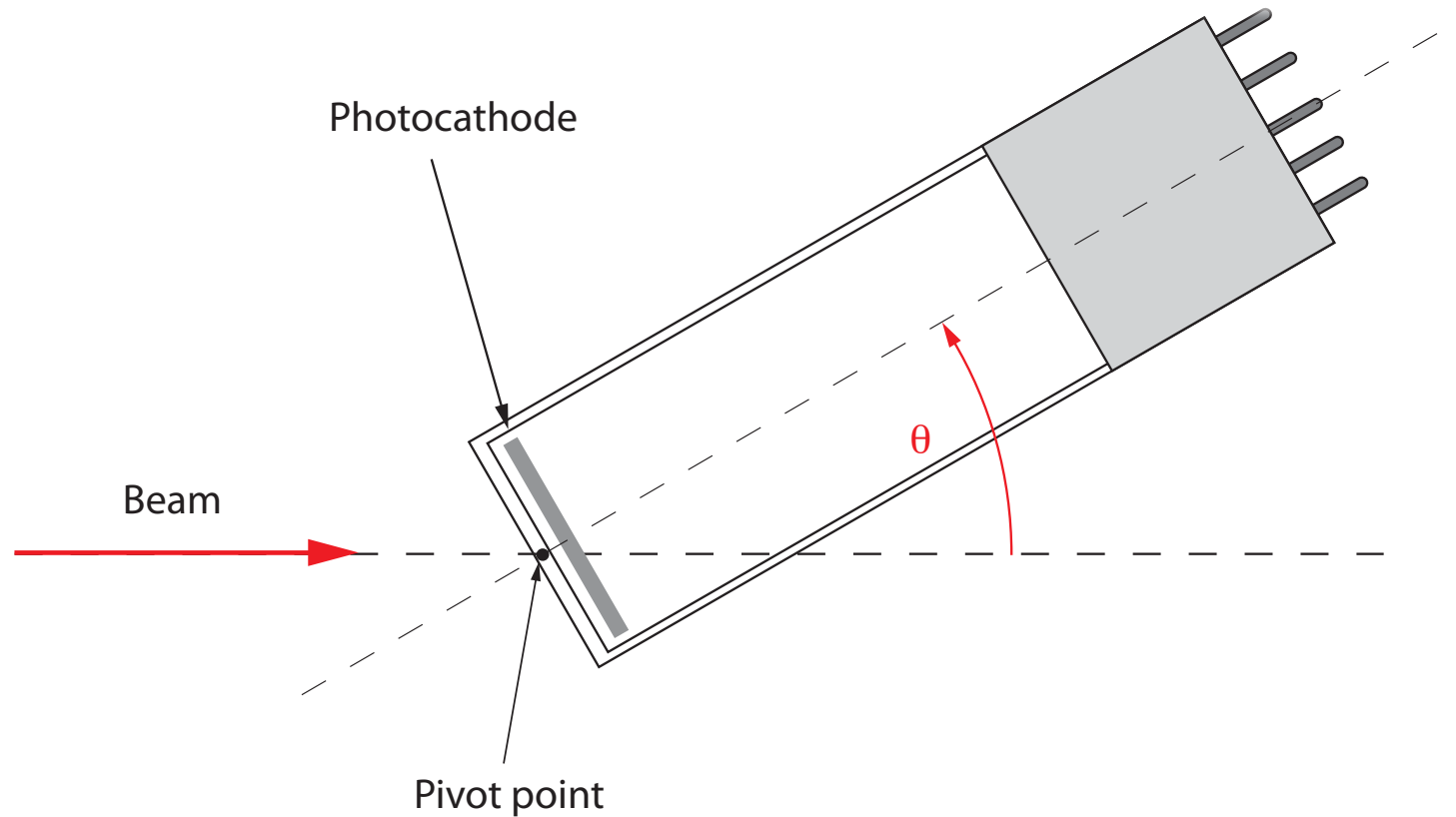


Test beam data show that a through-going muon generates ~ 30 *pe* in a PMT window (R7525), equivalent to ~ 120 GeV signal in this particular case, due to the Cherenkov light produced in the window.

The average plano-concave window thickness is ~ 6 mm for R7525.

The following studies show that we can effectively suppress this effect by altering the window+photocathode interface.

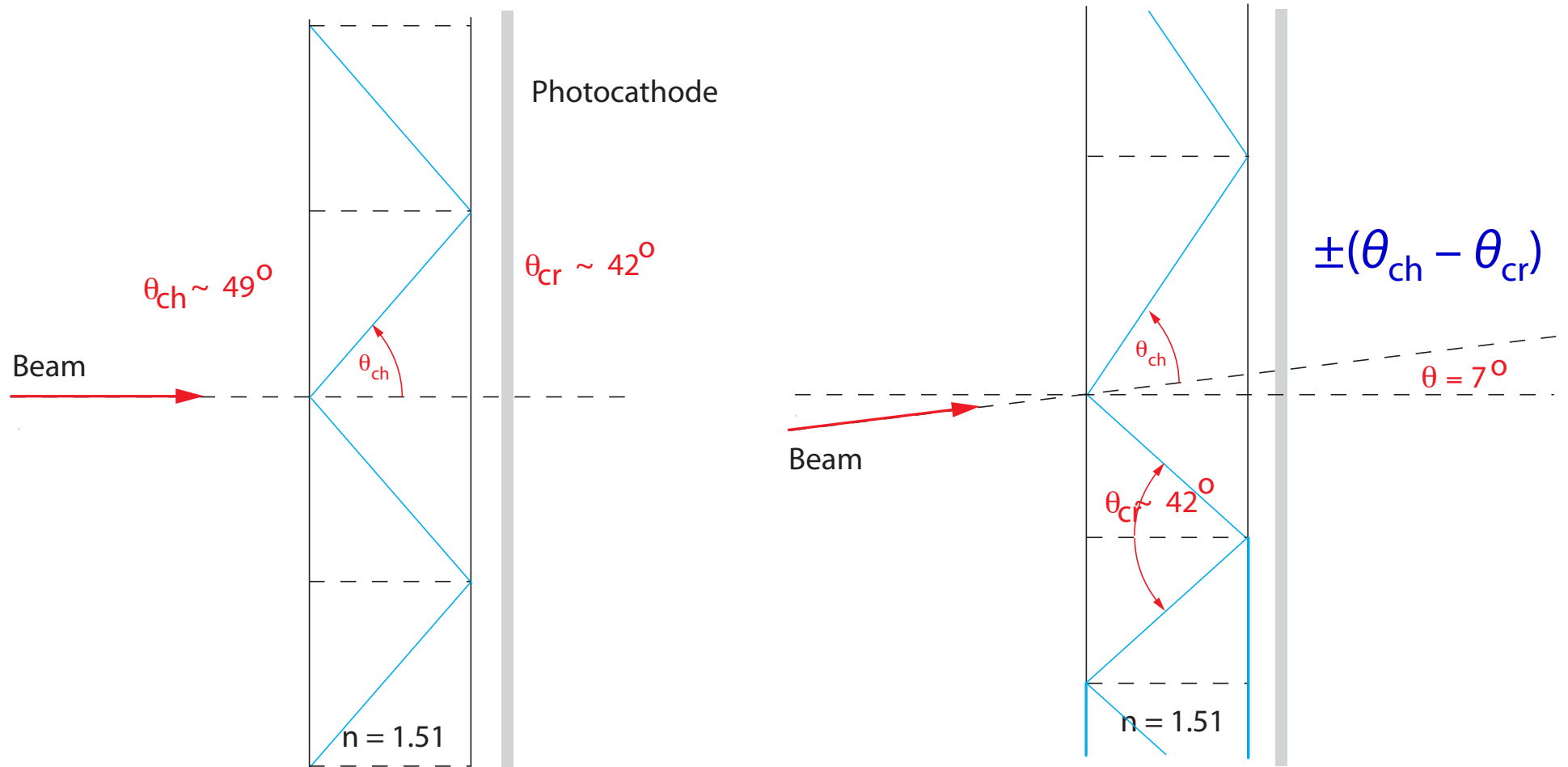
Idea



If the photocathode is separated from the window inside the PMT and deposited on a thin mica disk, then the pe emission will cease while the Cherenkov light is trapped in the glass

The critical measurement is the response of the PMT to relativistic charged particles as a function of the incidence angle, θ . We investigate this response under different conditions for future applications in calorimetry

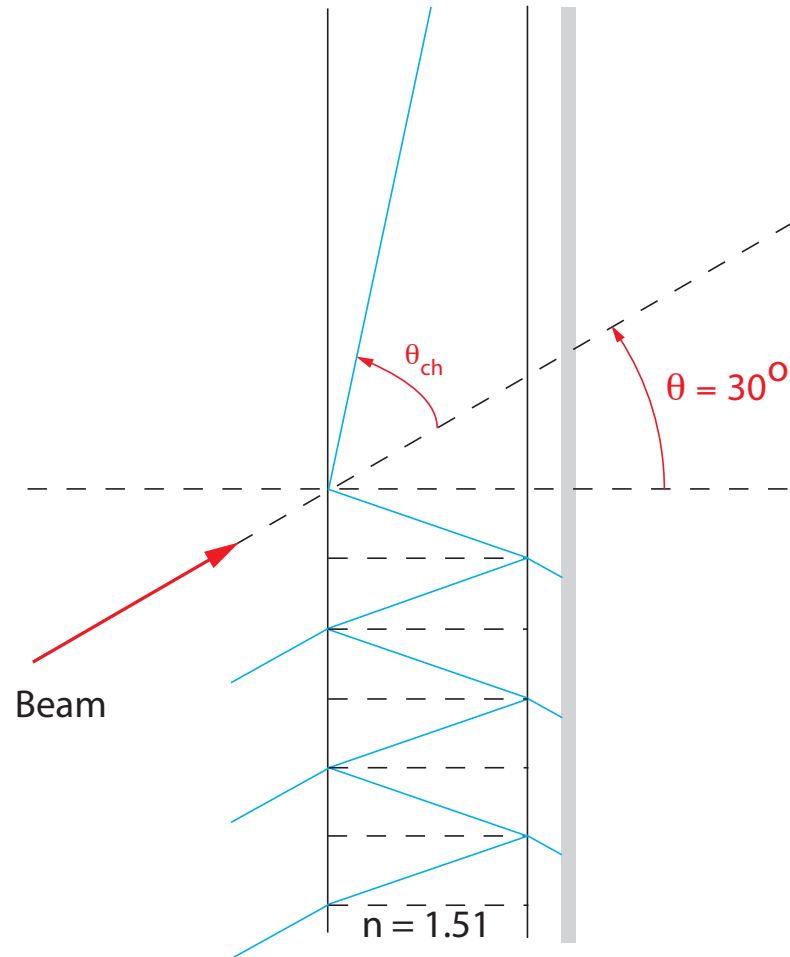
Capturing Cherenkov Light - I



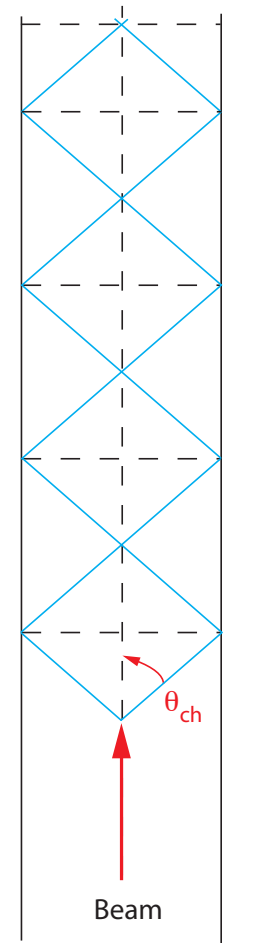
Case #1 (0°)

Case #2 (7°)

Capturing Cherenkov Light - II



Case #3 (30°)



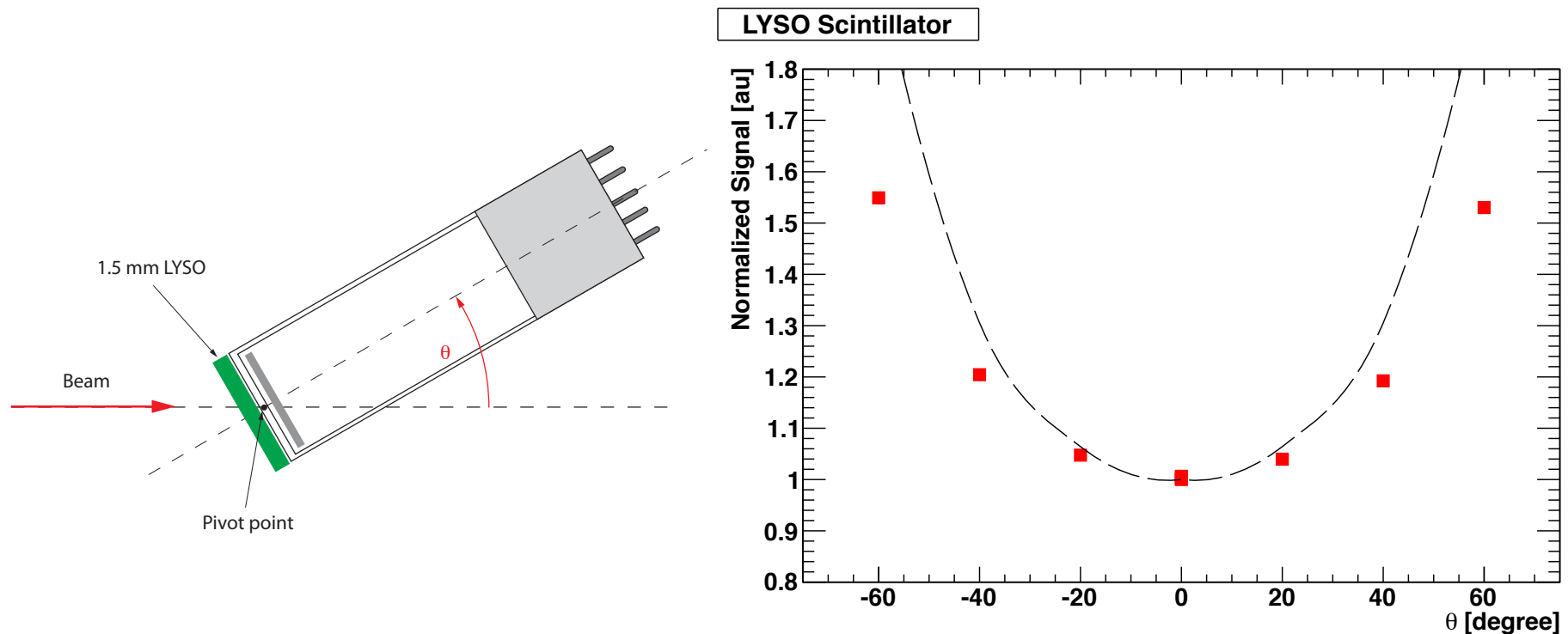
Case #4 (90°)

AntiCherenkov PMT (ET 865) Prototype

ET 865	Typical
Photocathode	Bialkali
Active diameter	23 mm
QE at peak	20 %
Luminous sensitivity	50 $\mu\text{A}/\text{lm}$
Dynodes	LF 8 stages
Dynode material	SbCs
Max anode sensitivity	50 A/lm
HV for max an. sens.	1200 V
Gain at nominal an.sens.	2×10^5
Dark current at max an. sens.	0.25 nA
-5% nonlinearity (A)	10 mA
Single <i>pe</i> peak/valley	1.8



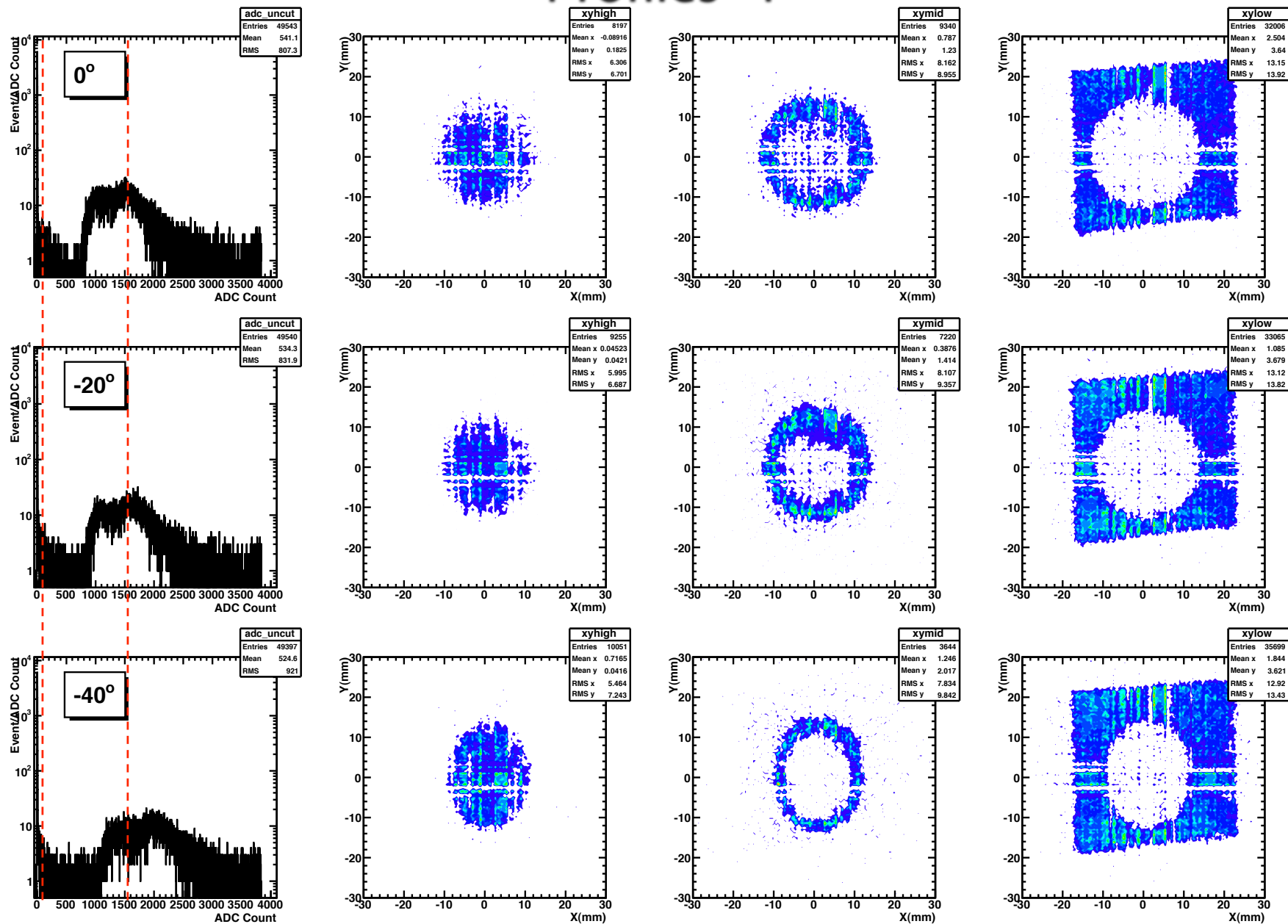
Angular Scan with 1.5 mm Thick Scintillator



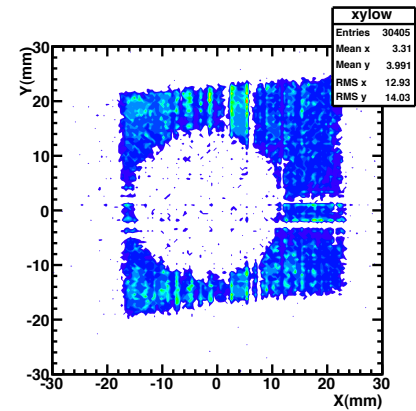
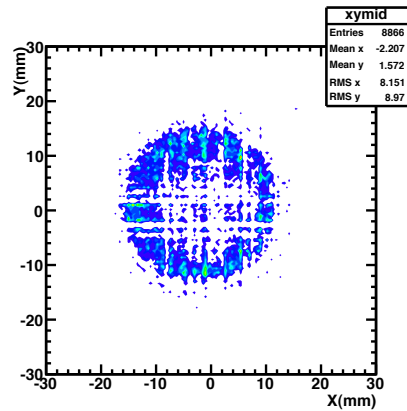
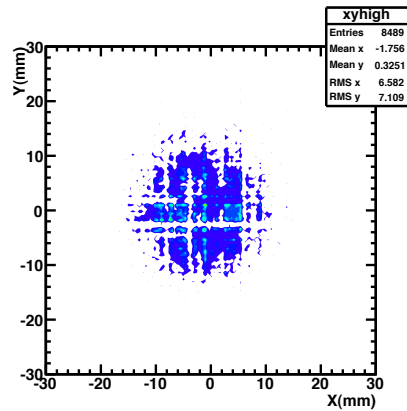
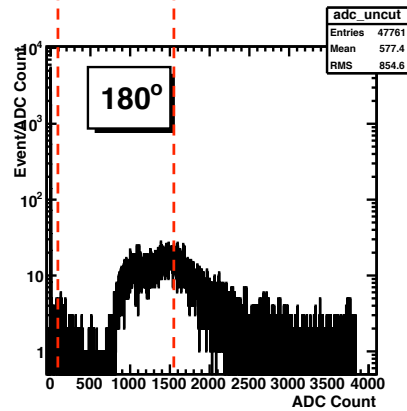
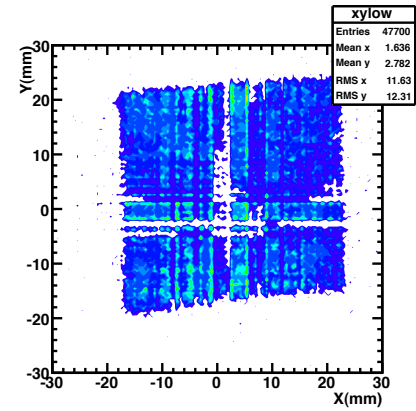
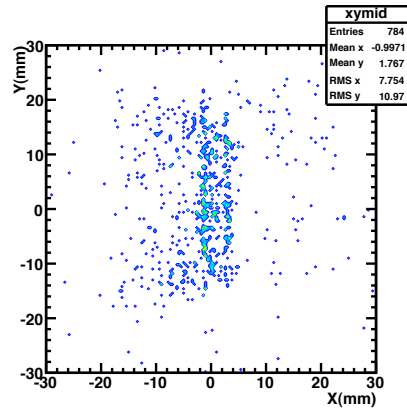
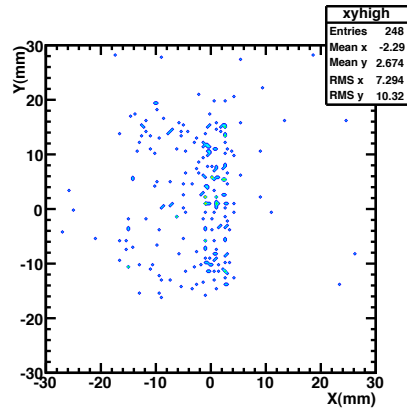
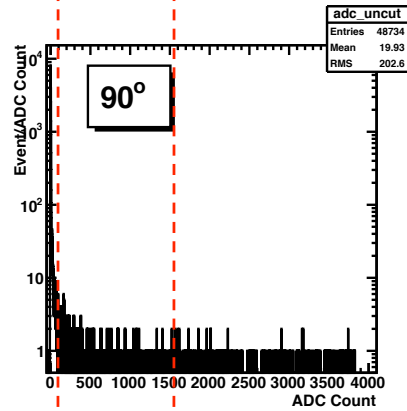
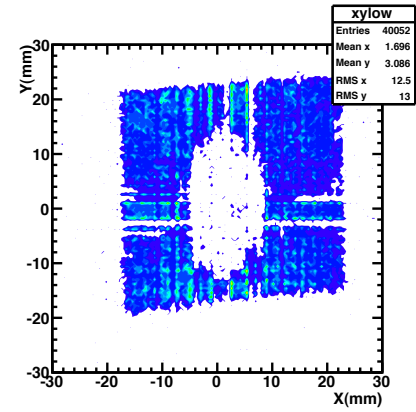
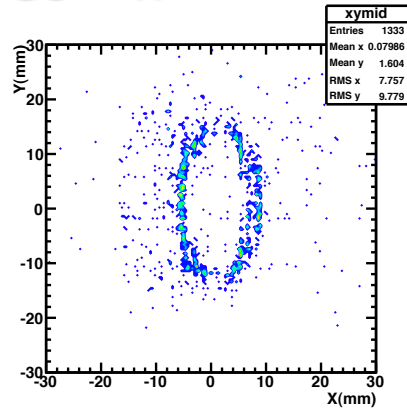
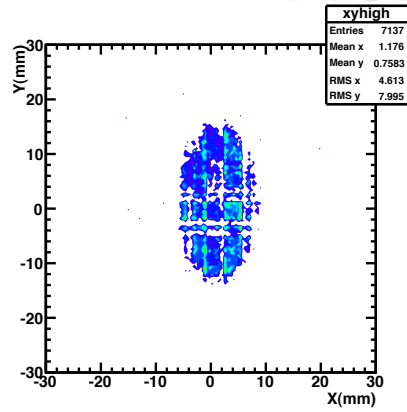
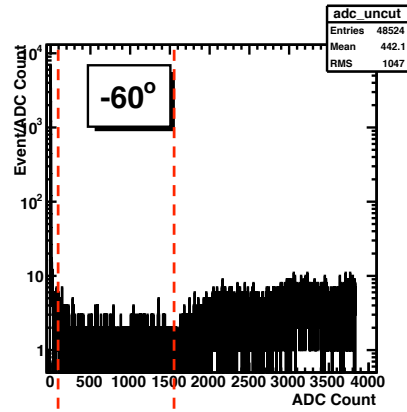
As a test, a 1.5-mm-thick LYSO scintillator disk is mounted on the AntiCherenkov PMT (1-1/8" dia.) and rotated around the pivot point. 180 GeV/c μ^+ beam is used and the impact point of the particles is measured by a beam chamber.

PMT is operated at -1100 V and $g \sim 10^5$.

Profiles - I

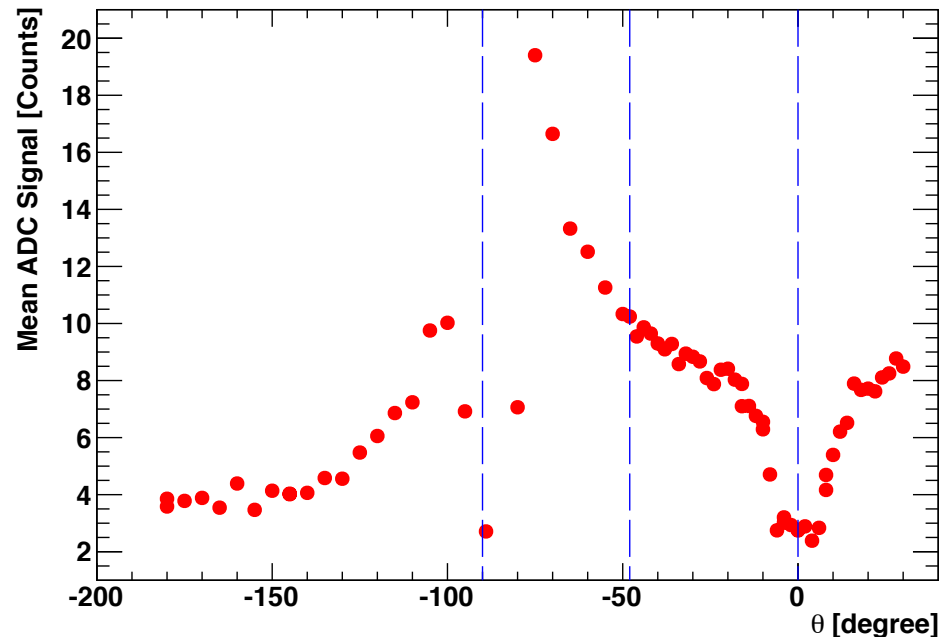


Profiles - II

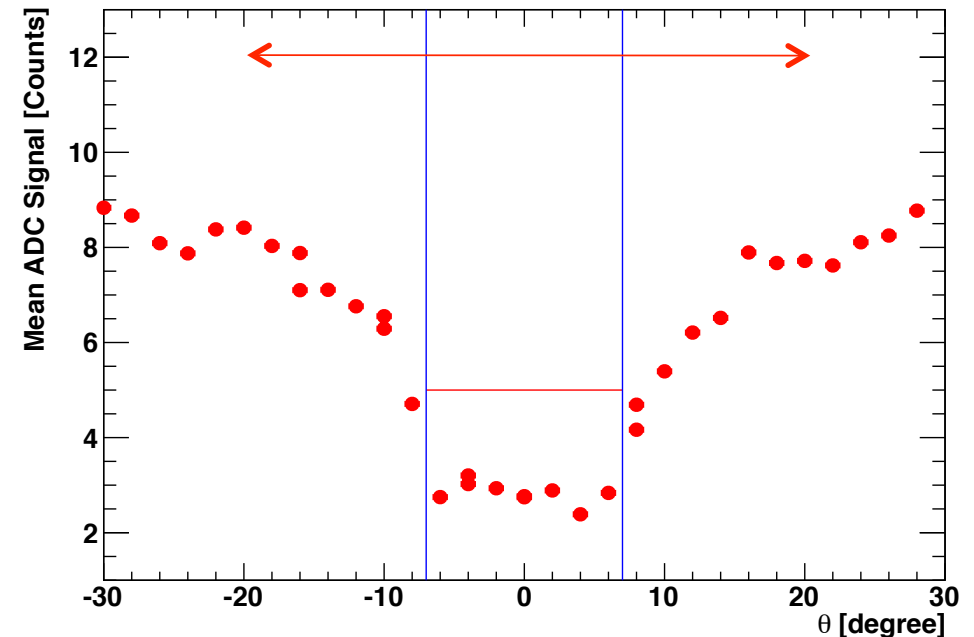


Angular Scan with Bare AntiCherenkov PMT

AntiCherenkov PMT 865B



AntiCherenkov PMT 865B

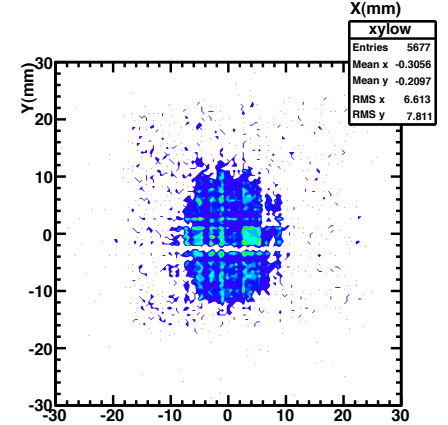
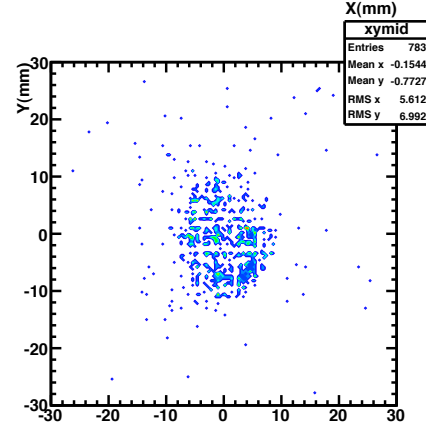
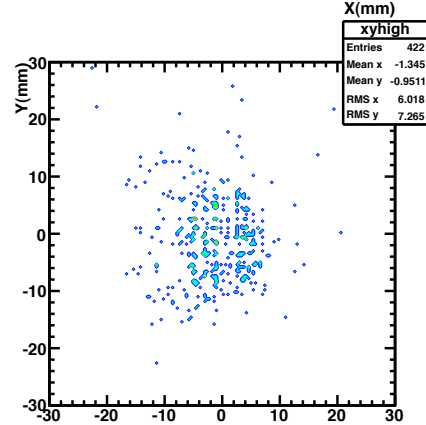
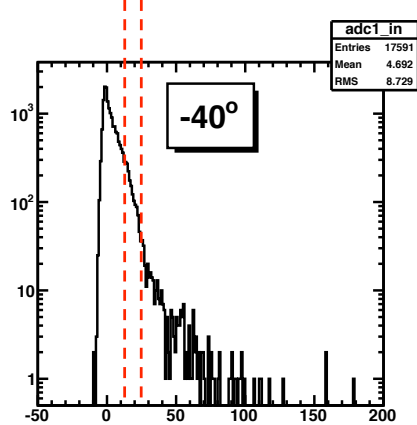
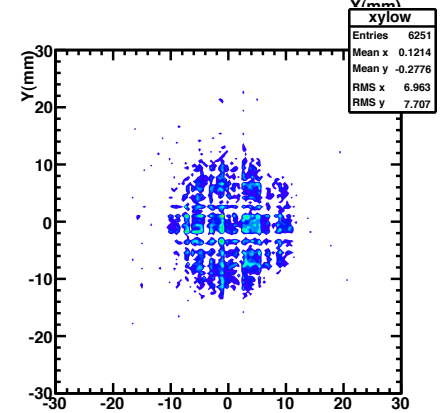
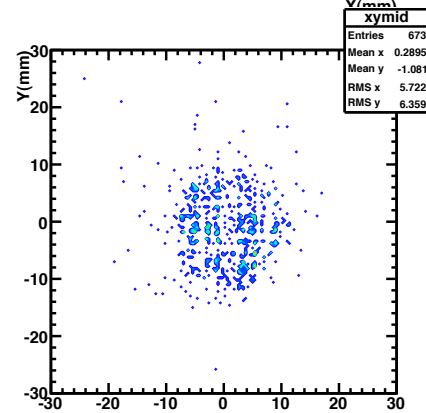
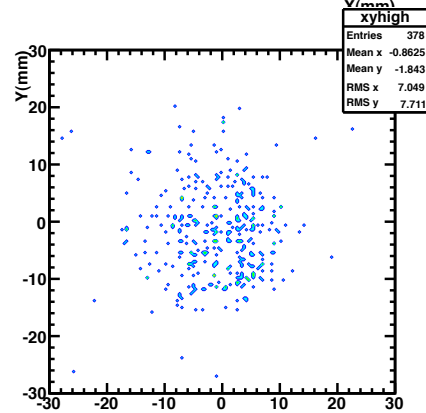
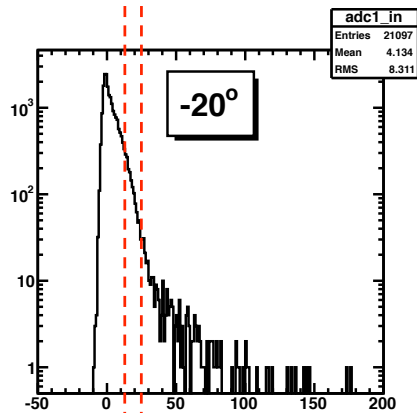
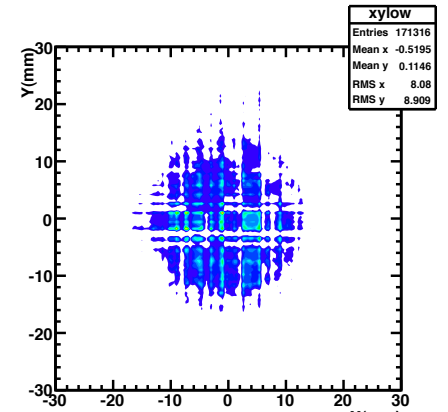
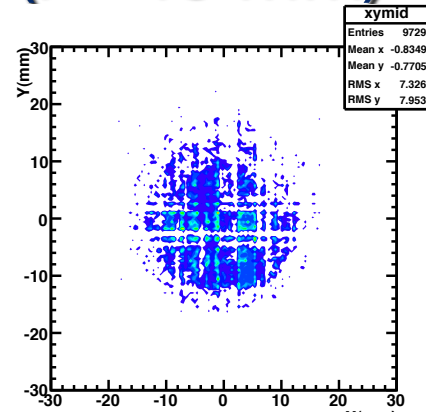
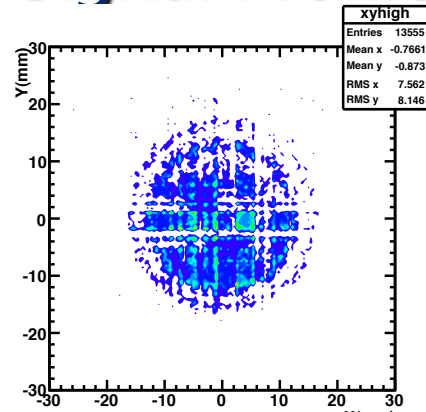
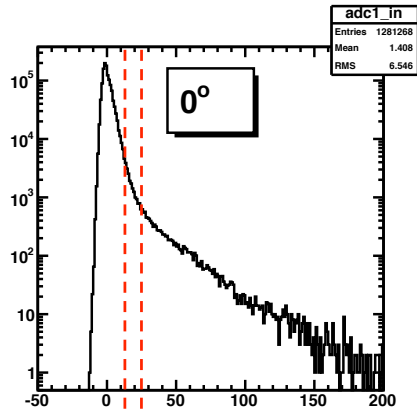


The angular scan of the bare PMT is carried out with 180 GeV/c μ^+ beam. A beam spot of 10 mm diameter is selected. Around zero, $\pm(\theta_{\text{ch}} - \theta_{\text{cr}}) = \pm 7^\circ$, and 90° , the Cherenkov light is trapped in the window glass.

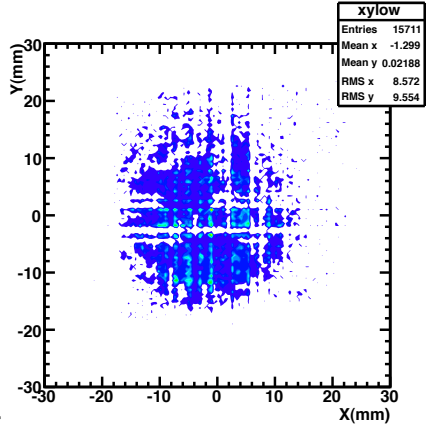
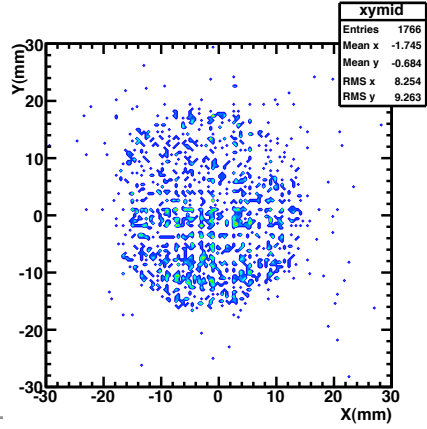
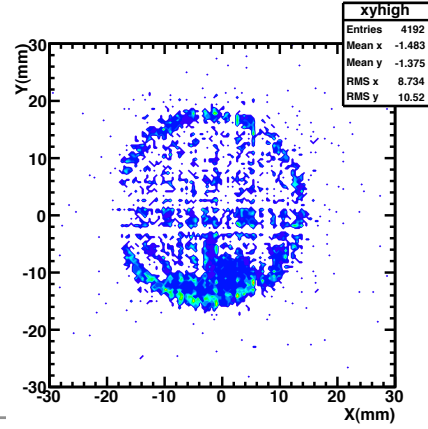
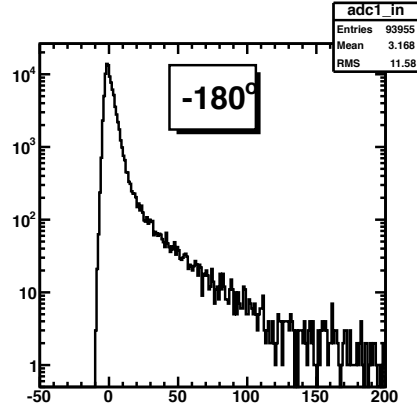
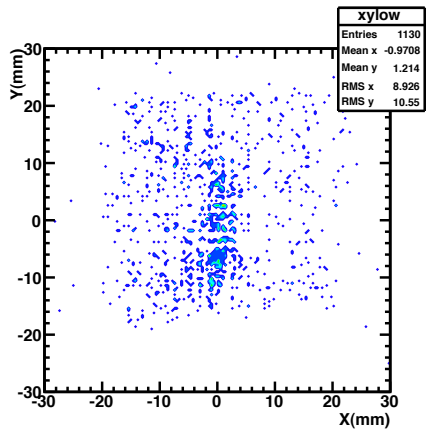
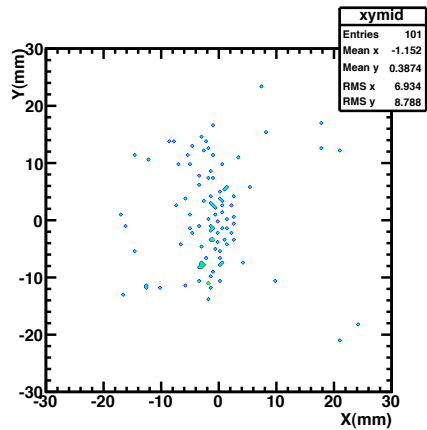
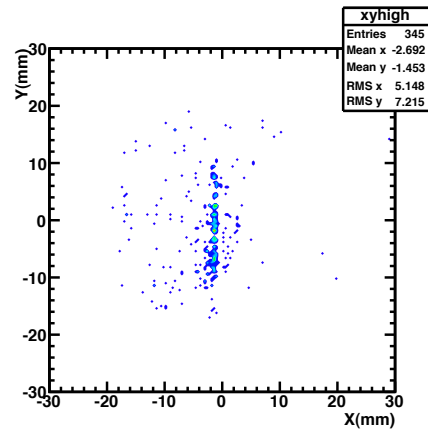
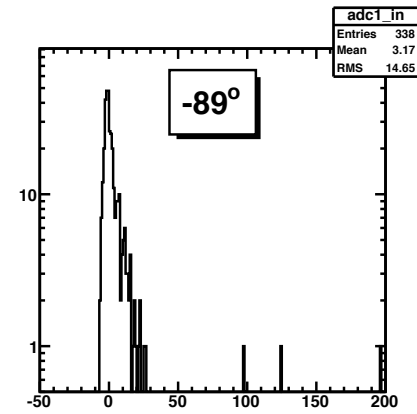
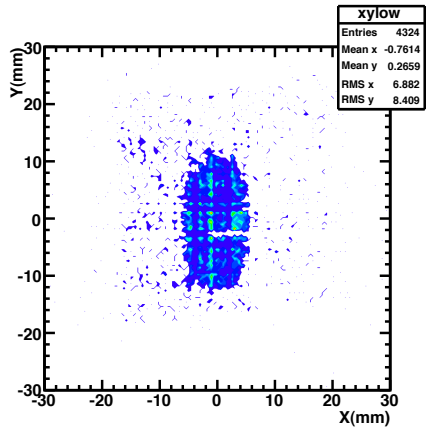
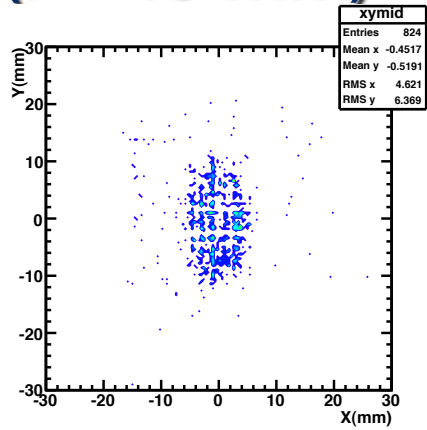
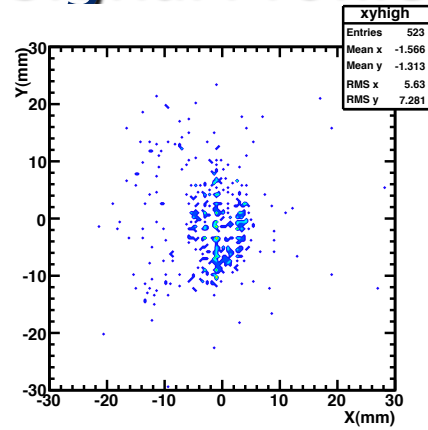
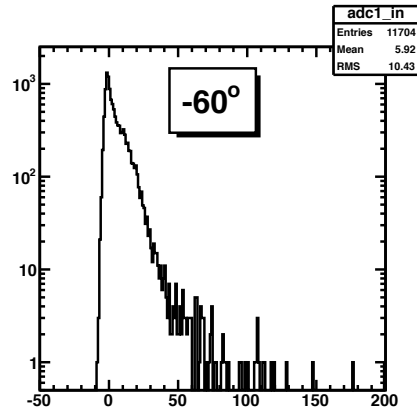
The effectiveness is maintained up to large angles ($\sim 40^\circ$). At larger angles ($90^\circ < \theta < \theta_{\text{ch}}$), the signal amplitude is $\propto 1/\cos(\theta)$.

The signal is suppressed at $\sim 90^\circ$ and $\sim 180^\circ$ for the same reasons.

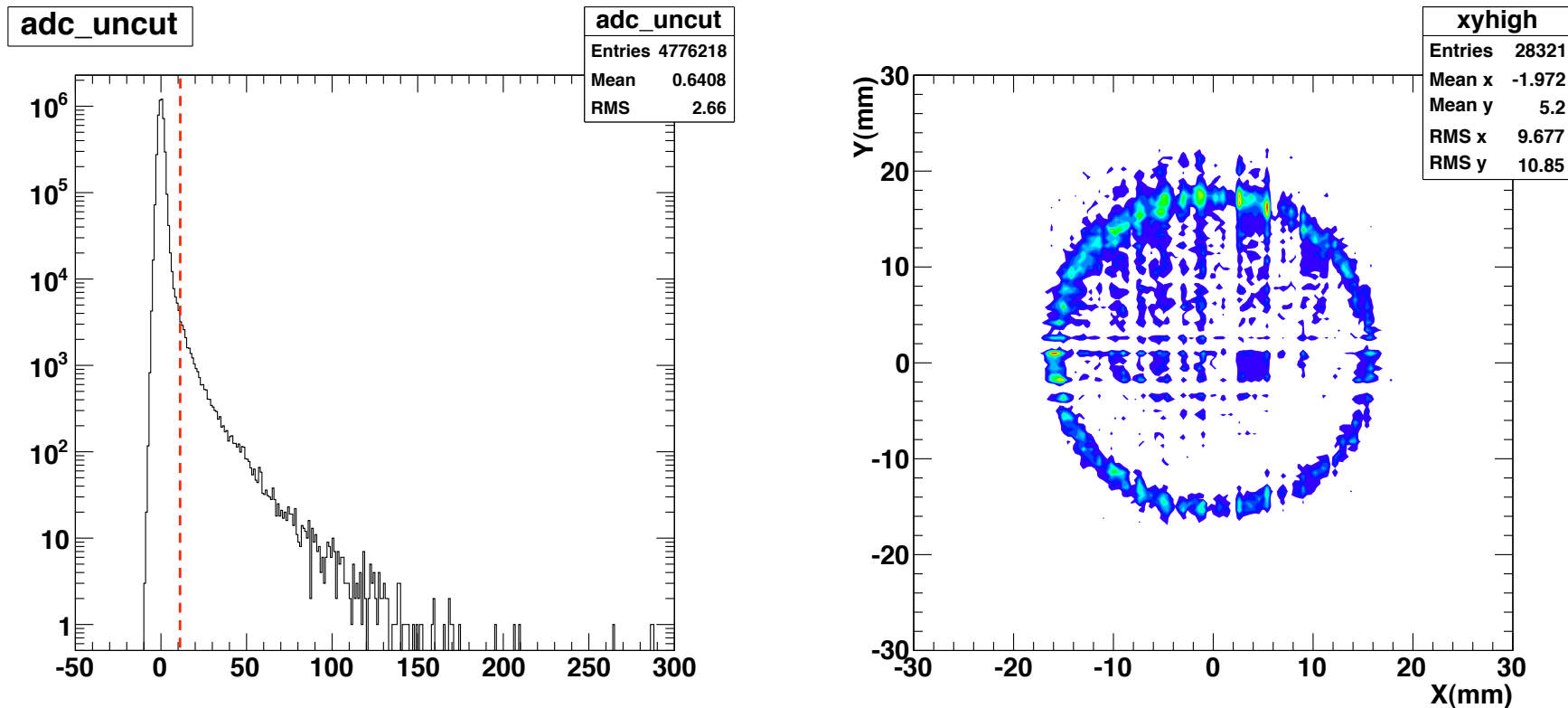
Signal Profiles ($r = 15$ mm) - I



Signal Profiles ($r = 15$ mm) - II

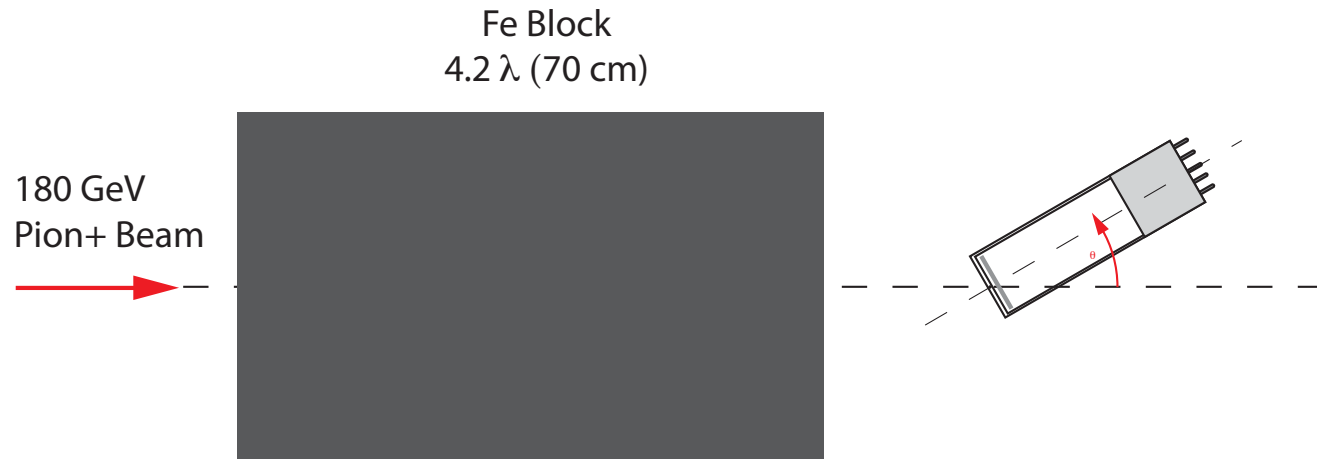


PMT Edge Effect

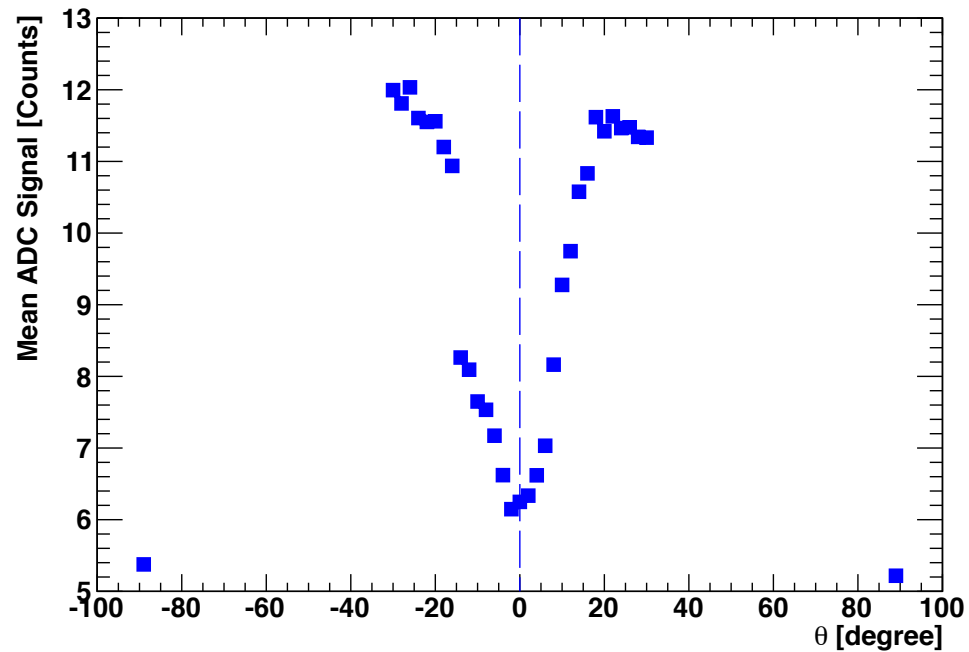


The largest contribution to the tail of the spectrum comes from the PMT rim when probed by muons. The central photocathode disk works as intended. The Cherenkov radiation from the disk ~negligible.

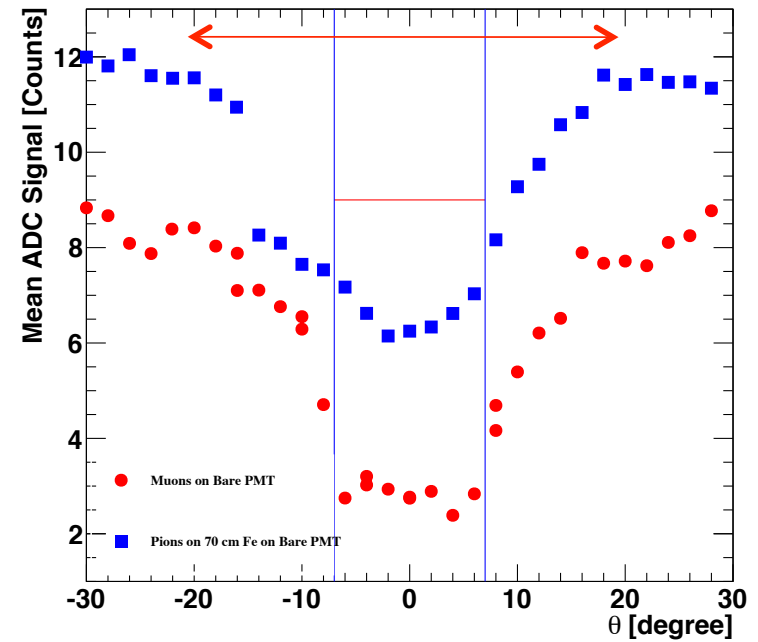
AntiCherenkov PMT behind a Calorimeter Absorber



AntiCherenkov PMT 865B 180 GeV π^+ on 70 cm Fe



AntiCherenkov PMT 865B 180 GeV π^+ on 70 cm Fe



Conclusions

AntiCherenkov PMT is effective suppressing signal due to the Cherenkov radiation produced in the window by relativistic charged particles

The most effective zone is around 0° incidence angle $\pm(\theta_{\text{ch}} - \theta_{\text{cr}})$. It is uniquely determined by the refractive index (n) of window glass ($n \sim 1.51$ and $\theta_{\text{ch}} - \theta_{\text{cr}} = 7^\circ$ in our case)

Larger n values will increase this region, *e.g.* sapphire will give $\pm 14^\circ$

It may be possible to “shape” the window to further increase the effective zone

AntiCherenkov PMTs are ideally suited for use behind calorimeters (late developing showers, punch-through muons *etc.*)

Newer high QE photocathodes will further enhance the capabilities of this PMT and compensate for the additional refractive surfaces (2) introduced by the displacement of the photocathode

Thanks go to Walter Kononenko, Andy Cormack, Paul Davisson, Larry Ludlum, Ron McAlpine who have paved the road way before us